

Statistical Physics & Condensed Matter Theory I: Exercise

A (classical) field theory for the Ising model

a) Trivial since $\sum_{S_a} e^{-\beta h_a S_a} = 2 \cosh h_a$.

b) We here consider the given representation of unity, but for a shifted value of the field:

$$\mathbf{1} \equiv \mathcal{N} \int \mathcal{D}\psi e^{-\frac{1}{4} \sum_{ab} (\psi_a + \theta_a)(K^{-1})_{ab} (\psi_b + \theta_b)}, \quad \mathcal{D}\psi \equiv \prod_{\mathbf{a}} d\psi_{\mathbf{a}}$$

in which $\theta_{\mathbf{a}}$ are constants to be chosen (the integration measure doesn't change under such shifts by constants). Written out explicitly,

$$\begin{aligned} \mathbf{1} &\equiv \mathcal{N} \int \mathcal{D}\psi e^{-\frac{1}{4} \sum_{ab} \psi_a (K^{-1})_{ab} \psi_b - \frac{1}{2} \sum_{ab} \theta_a (K^{-1})_{ab} \psi_b - \frac{1}{4} \sum_{ab} \theta_a (K^{-1})_{ab} \theta_b} \\ &= e^{-\frac{1}{4} \sum_{ab} \theta_a (K^{-1})_{ab} \theta_b} \times \mathcal{N} \int \mathcal{D}\psi e^{-\frac{1}{4} \sum_{ab} \psi_a (K^{-1})_{ab} \psi_b - \frac{1}{2} \sum_{ab} \theta_a (K^{-1})_{ab} \psi_b} \end{aligned}$$

Choosing $\theta_{\mathbf{a}} = -2 \sum_{\mathbf{b}} K_{ab} S_{\mathbf{b}}$ gives

$$e^{\sum_{ab} S_a K_{ab} S_b} = \mathcal{N} \int \mathcal{D}\psi e^{-\frac{1}{4} \sum_{ab} \psi_a (K^{-1})_{ab} \psi_b + \sum_{\mathbf{a}} S_{\mathbf{a}} \psi_{\mathbf{a}}}$$

which gives the required answer by direct substitution.

c) Simple.

d) Consider

$$x_{\mathbf{a}} = 2 \sum_{\mathbf{b}} K_{ab} \phi_{\mathbf{b}} = \frac{2}{N^{3/2}} \sum_{\mathbf{k}, \mathbf{k}'} \sum_{\mathbf{b}} e^{-i\mathbf{k} \cdot (\mathbf{a}-\mathbf{b}) - i\mathbf{k}' \cdot \mathbf{b}} K_{\mathbf{k}} \phi_{\mathbf{k}'} = \frac{2}{\sqrt{N}} \sum_{\mathbf{k}} e^{-i\mathbf{k} \cdot \mathbf{a}} K_{\mathbf{k}} \phi_{\mathbf{k}} \quad (1)$$

We can then write

$$\sum_{\mathbf{a}} \frac{x_{\mathbf{a}}^2}{2} = \frac{2}{N} \sum_{\mathbf{k}_1, \mathbf{k}_2} K_{\mathbf{k}_1} K_{\mathbf{k}_2} \phi_{\mathbf{k}_1} \phi_{\mathbf{k}_2} \sum_{\mathbf{a}} e^{-i(\mathbf{k}_1 + \mathbf{k}_2) \cdot \mathbf{a}} = 2 \sum_{\mathbf{k}} K_{\mathbf{k}} K_{-\mathbf{k}} \phi_{\mathbf{k}} \phi_{-\mathbf{k}} \quad (2)$$

and

$$\sum_{\mathbf{a}} \frac{x_{\mathbf{a}}^4}{12} = \frac{4}{3N^2} \sum_{\mathbf{k}_1, \dots, \mathbf{k}_4} K_{\mathbf{k}_1} K_{\mathbf{k}_2} K_{\mathbf{k}_3} K_{\mathbf{k}_4} \phi_{\mathbf{k}_1} \phi_{\mathbf{k}_2} \phi_{\mathbf{k}_3} \phi_{\mathbf{k}_4} \sum_{\mathbf{a}} e^{-i(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4) \cdot \mathbf{a}} \quad (3)$$

Doing the gradient expansion in each of these, and keeping only leading terms,

$$\sum_{\mathbf{a}} \frac{x_{\mathbf{a}}^2}{2} \simeq 2 \sum_{\mathbf{k}} (K_0 + \frac{1}{2} |\mathbf{k}|^2 K_0'' + \dots)^2 \phi_{\mathbf{k}} \phi_{-\mathbf{k}} \quad (4)$$

$$\sum_{\mathbf{a}} \frac{x_{\mathbf{a}}^4}{12} = \frac{4K_0^4}{3N} \sum_{\mathbf{k}_1 \dots \mathbf{k}_4} \phi_{\mathbf{k}_1} \phi_{\mathbf{k}_2} \phi_{\mathbf{k}_3} \phi_{\mathbf{k}_4} \delta_{\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4, 0} \quad (5)$$

Putting things together, we get the form for the action given in the question, with coefficients

$$c_1 = K_0 - 2K_0^2, \quad c_2 = \frac{1}{2}K_0'' - 2K_0K_0'', \quad c_3 = -1, \quad c_4 = \frac{4}{3}K_0^4 \quad (6)$$

e) Back to real space (specializing to $h = 0$)

$$S[\phi] = \int d^d x [c_2(\partial\phi)^2 + c_1\phi^2 + c_4\phi^4] \quad (7)$$

Rescaling $\phi \rightarrow \frac{1}{\sqrt{2c_2}}\phi$, we get $r = \frac{c_1}{c_2}$ and $g = \frac{c_4}{4c_2^2}$.

The calculation only makes sense if $c_2 > 0$ so $K_0 < 1/4$. r changes sign when $c_1 = 0$ so $K_0 = 1/2$ and thus the critical temperature is given by the condition $\beta_c = \frac{1}{2J_0}$. For $\beta > \beta_c$, $r < 0$ and the system is in the ordered phase. For $\beta < \beta_c$, $r > 0$ and the system is in the disordered phase.